

Quantum shift

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Quantum partial derivatives and argument shift method

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Outline

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Definition (Algebra)

An *algebra* is a vector space A with an associative bilinear mapping

$$A^2 \rightarrow A, \quad (x, y) \mapsto xy.$$

- The set $M(d, \mathbb{C})$ of d by d complex matrices is a complex algebra.
- Suppose that H is a Hilbert space. The set $B(H)$ of bounded linear operators on the Hilbert space H is a complex algebra.

Commutator and Lie algebra

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Remark (Commutator)

The alternating mapping

$$A^2 \rightarrow A, \quad (x, y) \mapsto [x, y] = xy - yx$$

satisfies the Jacobi identity.

Definition (Lie algebra)

A *Lie algebra* is a vector space g with an alternating mapping

$$g^2 \rightarrow g, \quad (x, y) \mapsto [x, y]$$

satisfying the Jacobi identity.

General linear Lie algebra

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- 1 Any algebra is a Lie algebra.
- 2 Any Lie algebra is isomorphic to a Lie subalgebra of some algebra by the Poincaré–Birkhoff–Witt theorem.

Remark

The notion of Lie algebras is an abstraction of commutators.

Definition (General linear Lie algebra)

The general linear Lie algebra $gl(d, \mathbb{C}) = M(d, \mathbb{C})$.

Ado's theorem

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Theorem (Ado)

Any finite dimensional complex Lie algebra g is isomorphic to a Lie subalgebra of the general linear Lie algebra $gl(d, \mathbb{C})$ for some d .

Suppose that (e_1, \dots, e_d) is a basis of a finite dimensional complex Lie algebra g .

Symmetric algebra

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Definition (Symmetric algebra)

The symmetric algebra Sg is an algebra generated by the elements (e_1, \dots, e_d) subject to the defining relations

$$e_m e_n = e_n e_m, \quad m, n = 1, \dots, d.$$

Remark

The symmetric algebra Sg is nothing but the polynomial algebra $\mathbb{C}[e_1, \dots, e_d]$.

Universal enveloping algebra

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Definition (Universal enveloping algebra)

The universal enveloping algebra Ug is an algebra generated by the elements (e_1, \dots, e_d) subject to the defining relations

$$e_m e_n = e_n e_m + [e_m, e_n], \quad m, n = 1, \dots, d.$$

Remark

The universal enveloping algebra Ug is a quantization of the symmetric algebra Sg .

Lie–Poisson bracket

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Theorem (Lie–Poisson bracket)

There exists a unique Poisson bracket on the symmetric algebra Sg extending the Lie bracket on the Lie algebra g .

$$\begin{array}{ccc} Sg \times Sg & \xrightarrow{\text{Poisson bracket}} & Sg \\ \uparrow & & \uparrow \\ g \times g & \xrightarrow{\text{Lie bracket}} & g \end{array}$$

Remark

We have $\text{gr } Ug = Sg$.

Classical argument shift method

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Suppose that ξ is an element of the dual space g^* and let

$$\bar{\partial}_\xi = \xi(e_1) \frac{\partial}{\partial e_1} + \cdots + \xi(e_d) \frac{\partial}{\partial e_d}$$

be a derivation on the symmetric algebra Sg .

Theorem (Mishchenko and Fomenko, 1978)

The family

$$\bigcup_{n=0}^{\infty} \left\{ \bar{\partial}_\xi^n x : x \text{ is Poisson central} \right\} \quad (1)$$

is Poisson commutative.

Argument shift algebra (classical and quantum)

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Definition (Argument shift algebra)

The Poisson commutative subalgebra \overline{C}_ξ generated by the Poisson commutative family (1) is called the argument shift algebra in the direction ξ .

Definition (Quantum argument shift algebra)

A commutative subalgebra C_ξ of the universal enveloping algebra Ug satisfying $\text{gr } C_\xi = \overline{C}_\xi$ is called a quantum argument shift algebra in the direction ξ .

The motivation for my talk is to quantize not only the algebra \overline{C}_ξ but also the derivation $\overline{\partial}_\xi$.

Generating matrix of general linear Lie algebra

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Let $e = \begin{pmatrix} e_1^1 & \cdots & e_d^1 \\ \cdots & \cdots & \cdots \\ e_1^d & \cdots & e_d^d \end{pmatrix}$ be a matrix satisfying the following.

- The set

$$\left\{ e_j^i : i, j = 1, \dots, d \right\}$$

is a basis of the general linear Lie algebra $gl(d, \mathbb{C})$.

- We have the commutation relations

$$[e_{j_1}^{i_1}, e_{j_2}^{i_2}] = \delta_{j_2}^{i_1} e_{j_1}^{i_2} - \delta_{j_1}^{i_2} e_{j_2}^{i_1}.$$

Partial derivatives

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We define

$$\bar{\partial}x = \begin{pmatrix} \bar{\partial}_1^1 x & \dots & \bar{\partial}_d^1 x \\ \dots & \dots & \dots \\ \bar{\partial}_1^d x & \dots & \bar{\partial}_d^d x \end{pmatrix}, \quad \bar{\partial}_j^i = \frac{\partial}{\partial e_j^i}$$

for any element x of the symmetric algebra $Sgl(d, \mathbb{C})$.

Differential operator

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Remark

The differential operator

$$Sgl(d, \mathbb{C}) \rightarrow M(d, Sgl(d, \mathbb{C})), \quad x \mapsto \bar{\partial}x$$

is a unique linear operator satisfying the following.

- 1 $\bar{\partial}\nu = 0$ for any scalar ν .
- 2 $\bar{\partial}\text{tr}(\xi e) = \xi$ for any numerical matrix ξ .
- 3 (Leibniz rule)

$$\bar{\partial}(xy) = (\bar{\partial}x)y + x(\bar{\partial}y)$$

for any elements x and y of the symmetric algebra $Sgl(d, \mathbb{C})$.

Quantum differential operator

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Definition (Gurevich, Pyatov, and Saponov, 2012)

The quantum differential operator

$$Ugl(d, \mathbb{C}) \rightarrow M(d, Ugl(d, \mathbb{C})), \quad x \mapsto \partial x$$

is a unique linear operator satisfying the following.

- 1 $\partial \nu = 0$ for any scalar ν .
- 2 $\partial \operatorname{tr}(\xi e) = \xi$ for any numerical matrix ξ .
- 3 (quantum Leibniz rule)

$$\partial(xy) = (\partial x)y + x(\partial y) + (\partial x)(\partial y)$$

for any elements x and y of the universal enveloping algebra $Ugl(d, \mathbb{C})$.

Quantum argument shift method

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Suppose that ξ is a numerical matrix and let $\partial_\xi = \text{tr}(\xi\partial)$.

Theorem (I. and Sharygin, 2023)

The family

$$\bigcup_{n=0}^{\infty} \left\{ \partial_\xi^n x : x \text{ is central} \right\} \quad (2)$$

is commutative.

Corollary

The subalgebra C_ξ generated by the family (2) is a quantum argument shift algebra in the direction ξ .

Gelfand invariants

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Remark (Gelfand invariants)

The center C of the universal enveloping algebra $Ugl(d, \mathbb{C})$ is generated by the elements $\{\text{tr } e, \dots, \text{tr } e^d\}$.

Quantum differential operator (modified)

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Remark

The linear operator

$$Ugl(d, \mathbb{C}) \rightarrow M(d, Ugl(d, \mathbb{C})), \quad x \mapsto \text{diag}(x, \dots, x) + \partial x$$

is an algebraic homomorphism and **will be denoted by ∂** from now on. We have the *quantum Leibniz rule*

$$\partial(xy) = (\partial x)(\partial y)$$

for any elements x and y of the universal enveloping algebra $Ugl(d, \mathbb{C})$.

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We define

$$f_{\pm}^{(n)}(x) = \sum_{m=0}^n \frac{1 \pm (-1)^{n-m}}{2} \binom{n-1}{m} x^m.$$

Theorem (I, 2022)

We have

$$\begin{aligned} \partial(e^n)_j^i &= \sum_{m=0}^n (f_+^{(n-m)}(e)(e^m)_j^i + f_-^{(n-m)}(e)_j^i(e^m)^i) \\ &= \sum_{m=0}^n ((e^m)_j^i f_+^{(n-m)}(e) + (e^m)_j^i f_-^{(n-m)}(e)^i). \end{aligned}$$

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We assume the following form

$$\partial(e^n)_j^i = \sum_{m=0}^n (g_m^{(n)}(e)(e^m)_j^i + h_m^{(n)}(e)_j(e^m)^i),$$

where $g_m^{(n)}$ and $h_m^{(n)}$ are polynomials.

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We have

$$\begin{aligned}\partial(e^{n+1})_j^i &= \sum_{k=1}^d \partial((e^n)_k^i e_j^k) = \sum_{k=1}^d (\partial(e^n)_k^i) (\partial e_j^k) \\ &= \sum_{m=0}^n \sum_{k=1}^d (g_m^{(n)}(e)(e^m)_k^i + h_m^{(n)}(e)_k(e^m)^i) (e_j^k + E_j^k)\end{aligned}$$

by the quantum Leibniz rule and the induction hypothesis.

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We have

$$\begin{aligned}\partial(e^{n+1})_j^i &= \sum_{m=0}^n h_m^{(n)}(e) e^m \delta_j^i + \sum_{m=1}^{n+1} g_{m-1}^{(n)}(e) (e^m)_j^i \\ &\quad + \sum_{m=0}^n (g_m^{(n)}(e) + h_m^{(n)}(e) e)_j (e^m)^i \\ &= \sum_{m=0}^{n+1} (g_m^{(n+1)}(e) (e^m)_j^i + h_m^{(n+1)}(e)_j (e^m)^i)\end{aligned}$$

by the commutation relations

$$[(e^m)^i, e_j^k] = (e^m)^k \delta_j^i - \delta^k (e^m)_j^i.$$

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We obtained the initial condition

$$g_0^{(0)}(x) = 1, \quad h_0^{(0)}(x) = 0$$

and the recursive relation

$$g_0^{(n+1)}(x) = \sum_{m=0}^n h_m^{(n)}(x)x^m,$$

$$g_m^{(n+1)}(x) = g_{m-1}^{(n)}(x), \quad 0 < m \leq n+1,$$

$$h_m^{(n+1)}(x) = g_m^{(n)}(x) + h_m^{(n)}(x)x, \quad 0 \leq m < n+1,$$

$$h_{n+1}^{(n+1)}(x) = 0.$$

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Its solution is

$$g_m^{(n)}(x) = f_+^{(n-m)}(x), \quad h_m^{(n)}(x) = f_-^{(n-m)}(x).$$

Quantum directional derivative

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We adhere to the convention that $\text{tr } e^{-1} = 1$. We have

$$\begin{aligned}\partial_\xi \prod_m \text{tr } e^{n_m} &= \sum_{m_1=-1}^{n_1} \text{tr } e^{m_1} \sum_{m_2=-1}^{n_2} \text{tr } e^{m_2} \cdots \text{tr} \left(\xi \prod_k f_+^{(n_k - m_k)}(e) \right), \\ \partial_\xi^2 \prod_m \text{tr } e^{n_m} &= \sum_{m_1=-1}^{n_1} \text{tr } e^{m_1} \sum_{m_2=-1}^{n_2} \text{tr } e^{m_2} \cdots \sum_{k_1=0}^{n_1 - m_1} \sum_{k_2=0}^{n_2 - m_2} \cdots \\ &\quad \text{tr} \left(\xi \prod_\ell f_+^{(k_\ell)}(e) \partial \text{tr} \left(\xi \prod_\ell f_+^{(n_\ell - m_\ell - k_\ell)}(e) \right) \right).\end{aligned}$$

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We define

$$C_\xi^{(0)} = C, \quad C_\xi^{(n)} = C_\xi^{(n-1)} [\partial_\xi^n C].$$

We have

$$\begin{aligned} C_\xi^{(1)} &= C_\xi^{(0)} \left[\text{tr}(\xi e^n) : n = 1, 2, \dots \right], \\ C_\xi^{(2)} &= C_\xi^{(1)} \left[(\tau_\xi \circ \sigma) \begin{pmatrix} 0 & P_n^T \\ P_m & 0 \end{pmatrix} : m, n = 0, 1, 2, \dots \right] \\ &= C_\xi^{(1)} \left[(\tau_\xi \circ \sigma) \begin{pmatrix} 0 & P_n^T \\ P_m & 0 \end{pmatrix} : |m - n| \leq 1 \right]. \end{aligned}$$

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We have

$$\frac{1}{2}(\tau_\xi \circ \sigma) \begin{pmatrix} 0 & P_n^T \\ P_n & 0 \end{pmatrix} = \sum_{m=0}^n \text{tr}(\xi e^m \xi e^n f_-^{(n-m)}(e)),$$

$$\begin{aligned} (\tau_\xi \circ \sigma) \begin{pmatrix} 0 & P_n^T \\ P_{n+1} & 0 \end{pmatrix} &= \sum_{m=0}^n \text{tr}(\xi e^m \xi (e^{n+1} f_-^{(n-m)}(e) \\ &\quad + e^n f_-^{(n-m+1)}(e))). \end{aligned}$$

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The generators are $\text{tr}(\xi e)$, $\text{tr}(\xi e^2)$, ... and

$$\text{tr}(\xi^2 e),$$

$$\text{tr}(2\xi^2 e^2 + \xi e \xi e),$$

$$\text{tr}(\xi^2 e^3 + \xi e \xi e^2),$$

$$\text{tr}(2\xi^2 e^4 + 2\xi e \xi e^3 + \xi e^2 \xi e^2 + \xi^2 e^2),$$

$$\text{tr}(\xi^2 e^5 + \xi e \xi e^4 + \xi e^2 \xi e^3 + \xi^2 e^3),$$

$$\text{tr}(2\xi^2 e^6 + 2\xi e \xi e^5 + 2\xi e^2 \xi e^4 + \xi e^3 \xi e^3 + 4\xi^2 e^4 + \xi e \xi e^3),$$

$$\text{tr}(\xi^2 e^7 + \xi e \xi e^6 + \xi e^2 \xi e^5 + \xi e^3 \xi e^4 + 3\xi^2 e^5 + \xi e \xi e^4), \dots$$

They are mutually commutative.

Deformation quantization

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Suppose that M is a smooth manifold and let

$$C^\infty(M) \times C^\infty(M) \rightarrow C^\infty(M), \quad (x, y) \mapsto \{x, y\}$$

be a Poisson bracket. A *star product*

$$C^\infty(M) \times C^\infty(M) \rightarrow C^\infty(M)[[\nu]],$$
$$(x, y) \mapsto x \star y = xy + \sum_{n=1}^{\infty} B_n(x, y) \nu^n$$

is called a deformation quantization if it satisfies the following.

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- 1 B_1, B_2, \dots are bidifferential operators.
- 2 Associativity.
- 3 We have

$$B_1(x, y) - B_1(y, x) = \left[\frac{x \star y - y \star x}{\nu} \right]_{\nu=0} = \{x, y\}$$

for any smooth functions x and y .

Theorem (Kontsevich, 2003)

Any Poisson manifold has a deformation quantization.

Deformation quantization

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The dual space g^* of a Lie algebra g is a Poisson manifold.

Remark

The image of the restriction of the star product on the product $Sg \times Sg$ is contained in the polynomial algebra $(Sg)[\nu]$.

It makes sense to put $\nu = 1$ and obtain the star product on the symmetric algebra Sg .

Theorem

The universal enveloping algebra Ug is isomorphic to the symmetric algebra Sg with the star product.

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