

Constructing quantum argument-shift algebras via quantized shift operators

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Suppose that ξ is a $d \times d$ complex matrix and let

$$\partial_\xi = \text{tr}(\xi \partial) = \sum_{i,j=1}^d \xi_i^j \partial_j^i,$$

where $\partial_j^i \in \text{hom } U\mathfrak{gl}_d$ are the **quantum derivations** introduced by Gurevich, Pyatov, and Saponov. The main theorem is the following:

Theorem (I. and Sharygin, 2024)

*Suppose that x and y are **central** elements of $U\mathfrak{gl}_d$. Then*

$$[\partial_\xi^m x, \partial_\xi^n y] = 0$$

for any nonnegative integers m and n .

Symmetric and Universal Enveloping Algebras

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Suppose that $\mathfrak{g}$ is a Lie algebra.

- Both the symmetric algebra $S\mathfrak{g}$ and the universal enveloping algebra $U\mathfrak{g}$ are associative algebras.
- The symmetric algebra $S\mathfrak{g}$ (commutative) gives classical descriptions and the universal enveloping algebra $U\mathfrak{g}$ (non-commutative) gives quantum descriptions of the Lie algebra $\mathfrak{g}$.
- The symmetric algebra $S\mathfrak{g} = \text{gr } U\mathfrak{g}$ is the associated graded **Poisson algebra** of the universal enveloping algebra $U\mathfrak{g}$.

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- The symmetric algebra $S\mathfrak{g}$ is a subalgebra of the smooth functions algebra $C^\infty \mathfrak{g}^*$.
- There exists a unique Poisson bracket (Lie–Poisson bracket) on the smooth functions algebra $C^\infty \mathfrak{g}^*$ extending the Lie bracket.

$$\begin{array}{ccc} C^\infty \mathfrak{g}^* \times C^\infty \mathfrak{g}^* & \xrightarrow{\text{Lie–Poisson bracket}} & C^\infty \mathfrak{g}^* \\ \uparrow & & \uparrow \\ \mathfrak{g} \times \mathfrak{g} & \xrightarrow{\text{Lie bracket}} & \mathfrak{g} \end{array}$$

- The dual space $\mathfrak{g}^*$ of the Lie algebra $\mathfrak{g}$ is a **Poisson manifold**.

Deformation Quantization

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Consider a deformation quantization of $C^\infty \mathfrak{g}^*$.

Remark

$$\begin{array}{ccc} C^\infty \mathfrak{g}^* \times C^\infty \mathfrak{g}^* & \xrightarrow{\star} & (C^\infty \mathfrak{g}^*)[[\nu]] \\ \uparrow & & \uparrow \\ S\mathfrak{g} \times S\mathfrak{g} & \xrightarrow{\star} & (S\mathfrak{g})[\nu] \end{array}$$

- It makes sense to put $\nu = 1$ and obtain the star product on the symmetric algebra $S\mathfrak{g}$.
- The universal enveloping algebra $U\mathfrak{g}$ is isomorphic to the symmetric algebra $S\mathfrak{g}$ with the star product at $\nu = 1$.

Prior Research and Motivation

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Take a basis $(e_n)_{n=1}^d$ of $\mathfrak{g}$ and let

$$\bar{\partial}_\xi = \sum_{n=1}^d \xi(e_n) \frac{\partial}{\partial e_n} \in \text{der } S\mathfrak{g}$$

be the directional derivative along $\forall \xi \in \mathfrak{g}^*$. Let $\bar{\mathcal{C}}$ be the Poisson center of $S\mathfrak{g}$. The following theorem is referred to as the argument shift method.

Theorem (A. Mishchenko and A. Fomenko, 1978)

The subset $\left\{ \bar{\partial}_\xi^n x : (n, x) \in \mathbb{N} \times \bar{\mathcal{C}} \right\}$ is Poisson commutative.

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- We obtain a Poisson commutative subalgebra $\overline{C}_\xi$ generated by these elements $\overline{\partial}_\xi^n x$.
- Recall $\text{gr } U\mathfrak{g} = S\mathfrak{g}$.
- Vinberg asked if the argument shift algebra $\overline{C}_\xi$ can be quantised to a commutative subalgebra C_ξ of the universal enveloping algebra $U\mathfrak{g}$ in a way that $\text{gr } C_\xi = \overline{C}_\xi$.
- Such C_ξ is called a quantum argument shift algebra.

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- **Vinberg's problem** has been resolved in two ways:
 - Twisted Yangians: Nazarov–Olshanski.
 - Symmetrisation mapping: Tarasov.
- Also resolved using the Feigin–Frenkel center:
 - for regular elements ξ : Feigin et al. and Rybnikov.
 - for simple Lie algebras of types A and C: Futorny–Molev and Molev–Yakimova.

Motivation

The purpose of my talk is to quantize not only the algebra $\overline{\mathcal{C}}_\xi$ but also the **operator** $\overline{\partial}_\xi$.

Quantum Derivation

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Let $e = \begin{pmatrix} e_1^1 & \dots & e_d^1 \\ \dots & \dots & \dots \\ e_1^d & \dots & e_d^d \end{pmatrix}$ be a matrix satisfying the following.

- The set

$$\left\{ e_j^i : i, j = 1, \dots, d \right\}$$

is a basis of the general linear Lie algebra $\mathfrak{gl}(d, \mathbb{C})$.

- We have the commutation relations

$$[e_{j_1}^{i_1}, e_{j_2}^{i_2}] = \delta_{j_2}^{i_1} e_{j_1}^{i_2} - \delta_{j_1}^{i_2} e_{j_2}^{i_1}.$$

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We define

$$\bar{\partial}x = \begin{pmatrix} \bar{\partial}_1^1 x & \dots & \bar{\partial}_d^1 x \\ \dots & \dots & \dots \\ \bar{\partial}_1^d x & \dots & \bar{\partial}_d^d x \end{pmatrix}, \quad \bar{\partial}_j^i = \frac{\partial}{\partial e_i^j}$$

for any element x of the symmetric algebra $S\mathfrak{gl}(d, \mathbb{C})$.

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The derivation

$$\mathrm{Sgl}_d \rightarrow M(d, \mathrm{Sgl}_d), \quad x \mapsto \bar{\partial}x$$

is a unique linear mapping satisfying the following.

- 1 $\bar{\partial}\nu = 0$ for any scalar ν .
- 2 $\bar{\partial}\mathrm{tr}(\xi e) = \xi$ for any numerical matrix ξ .
- 3 (Leibniz rule)

$$\bar{\partial}(xy) = (\bar{\partial}x)y + x(\bar{\partial}y)$$

for any elements x and y of the symmetric algebra Sgl_d .

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There is no such mapping on $U\mathfrak{gl}_d$ because it is non-commutative.

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Definition (Gurevich, Pyatov, and Saponov, 2012)

The quantum derivation

$$U\mathfrak{gl}_d \rightarrow M(d, U\mathfrak{gl}_d), \quad x \mapsto \partial x$$

is a unique linear mapping satisfying the following.

- 1 $\partial\nu = 0$ for any scalar ν .
- 2 $\partial \operatorname{tr}(\xi e) = \xi$ for any numerical matrix ξ .
- 3 (quantum Leibniz rule)

$$\partial(xy) = (\partial x)y + x(\partial y) + (\partial x)(\partial y)$$

for any elements x and y of the universal enveloping algebra $U\mathfrak{gl}_d$.

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Let C be the center of the algebra $U\mathfrak{gl}_d$. Suppose that ξ is a numerical matrix and let $\partial_\xi = \text{tr}(\xi\partial)$. The main theorem is the following:

Theorem (I. and Sharygin, 2024)

The subset

$$\left\{ \partial_\xi^n x : (n, x) \in \mathbb{N} \times C \right\} \quad (1)$$

is commutative.

Corollary

The subalgebra C_ξ generated by the subset (1) is the quantum argument shift algebra in the direction ξ .

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- We may assume that $\xi = \text{diag}(z_1, \dots, z_d)$ is diagonal and $(z_1, \dots, z_d)$ is distinct considering the adjoint action of the general linear Lie group GL_d .
- Vinberg and Rybnikov showed that the quantum argument shift algebra in the direction ξ is the centralizer of the set

$$\left\{ e_i^j, \sum_{j \neq i} \frac{e_i^j e_j^i}{z_i - z_j} \right\}_{i=1}^d. \quad (2)$$

- Since, by definition, the quantum argument shift algebra is commutative, the proof is carried out by showing that the quantum argument shift $\partial_\xi^n x$ commutes with the elements (2) by induction on the natural number n .

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The center C of the algebra $U\mathfrak{gl}_d$ is the free commutative algebra on the elements

$$\mathrm{tr} e, \quad \dots, \quad \mathrm{tr} e^d.$$

They are called the Gelfand invariants. We would like to calculate the quantum argument shift $\partial_\xi^n x$ for a central element x . It is necessary and even sufficient to calculate the quantum derivation $\partial(e^n)_j^i$.

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Remark

The linear operator

$$U\mathfrak{gl}_d \rightarrow M(d, U\mathfrak{gl}_d), \quad x \mapsto \text{diag}(x, \dots, x) + \partial x$$

is an algebraic homomorphism and **will be denoted by ∂** from now on. We have the *quantum Leibniz rule*

$$\partial(xy) = (\partial x)(\partial y)$$

for any elements x and y of the universal enveloping algebra $U\mathfrak{gl}_d$.

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I obtained the following formula for the quantum derivation.

We define $f_{\pm}^{(n)}(x) = \sum_{m=0}^n \frac{1 \pm (-1)^{n-m}}{2} \binom{n-1}{m} x^m$.

Theorem (I, 2022)

We have

$$\begin{aligned} \partial(e^n)_j^i &= \sum_{m=0}^n (f_+^{(n-m)}(e)(e^m)_j^i + f_-^{(n-m)}(e)_j(e^m)^i) \\ &= \sum_{m=0}^n ((e^m)_j^i f_+^{(n-m)}(e) + (e^m)_j f_-^{(n-m)}(e)^i). \end{aligned}$$

The formula is used for the base case.

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We assume the following form

$$\partial(e^n)_j^i = \sum_{m=0}^n (g_m^{(n)}(e)(e^m)_j^i + h_m^{(n)}(e)_j(e^m)^i),$$

where $g_m^{(n)}$ and $h_m^{(n)}$ are polynomials. By the quantum Leibniz rule and the commutation relations

$$[(e^m)^i, e_j^k] = (e^m)^k \delta_j^i - \delta^k(e^m)_j^i,$$

We obtained the initial condition

$$g_0^{(0)}(x) = 1, \quad h_0^{(0)}(x) = 0$$

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and the recursive relation

$$g_0^{(n+1)}(x) = \sum_{m=0}^n h_m^{(n)}(x)x^m,$$

$$g_m^{(n+1)}(x) = g_{m-1}^{(n)}(x), \quad 0 < m \leq n+1,$$

$$h_m^{(n+1)}(x) = g_m^{(n)}(x) + h_m^{(n)}(x)x, \quad 0 \leq m < n+1,$$

$$h_{n+1}^{(n+1)}(x) = 0.$$

Its solution is

$$g_m^{(n)}(x) = f_+^{(n-m)}(x), \quad h_m^{(n)}(x) = f_-^{(n-m)}(x).$$

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The inductive step reduces to proving

$$[\mathrm{ad} e_i^j, \partial_\xi] = \left[\left[\mathrm{ad} \sum_{j \neq i} \frac{e_i^j e_j^i}{z_i - z_j}, \partial_\xi \right], \partial_\xi \right] = 0.$$

It can be done by computation.

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Our theorem causes a filtration

$$C_{\xi}^{(0)} = C, \quad C_{\xi}^{(n)} = C_{\xi}^{(n-1)} [\partial_{\xi}^n C]$$

of the quantum argument shift algebra C_{ξ} . Using the formula we obtain

$$C_{\xi}^{(1)} = C_{\xi}^{(0)} \left[\text{tr}(\xi e^n) : n = 1, 2, \dots \right],$$
$$C_{\xi}^{(2)} = C_{\xi}^{(1)} \left[\tau_{\xi} \left(\begin{pmatrix} 0 & P_n^{\top} \\ P_m & 0 \end{pmatrix} : m, n = 0, 1, 2, \dots \right) \right].$$

P_n : some matrix composed of binomial coefficients.

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But as for the second line these generators are redundant:

$$C_{\xi}^{(2)} = C_{\xi}^{(1)} \left[\tau_{\xi} \begin{pmatrix} 0 & P_n^{\top} \\ P_m & 0 \end{pmatrix} : |m - n| \leq 1 \right].$$

Lemma (I, 2025)

We have

$$\sigma \begin{pmatrix} 0 & P_m^{\top} \\ P_{m+2n} & 0 \end{pmatrix} = \sum_{k=0}^n \left(\binom{2n-k}{k} + \binom{2n-k-1}{k-1} \right) P_{m+k}^{(m+k)},$$
$$\sigma \begin{pmatrix} 0 & P_m^{\top} \\ P_{m+2n+1} & 0 \end{pmatrix} = \sum_{k=0}^n \binom{2n-k}{k} \left(P_{m+k+1}^{(m+k)} + P_{m+k}^{(m+k+1)} \right).$$

for any nonnegative integers m and n .

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Lemma reduces to the following relations.

1 For $\varepsilon = 0, 1$,

$$\binom{x+y+n}{2n+\varepsilon} + \binom{x-y+n}{2n+\varepsilon} = \sum_{m=0}^n \binom{x+m}{2m+\varepsilon} \left(\binom{y+n-m}{2(n-m)} + \binom{y-1+n-m}{2(n-m)} \right).$$

2
$$\sum_{m=0}^n \binom{x-m}{m} \binom{y+m}{n-m} = \sum_{m=0}^n \binom{x+y-m}{m} \binom{m}{n-m}.$$

3
$$\binom{x}{n} = \sum_{m=0}^n \binom{x-m}{m} \binom{m}{n-m} + \sum_{m=0}^{n-1} \binom{x-1-m}{m} \binom{m}{n-1-m}.$$

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They are shown by induction.

The generators are $\text{tr}(\xi e)$, $\text{tr}(\xi e^2)$, $\dots$ and

$$\text{tr}(\xi^2 e),$$

$$\text{tr}(2\xi^2 e^2 + \xi e \xi e),$$

$$\text{tr}(\xi^2 e^3 + \xi e \xi e^2),$$

$$\text{tr}(2\xi^2 e^4 + 2\xi e \xi e^3 + \xi e^2 \xi e^2 + \xi^2 e^2),$$

$$\text{tr}(\xi^2 e^5 + \xi e \xi e^4 + \xi e^2 \xi e^3 + \xi^2 e^3),$$

$$\text{tr}(2\xi^2 e^6 + 2\xi e \xi e^5 + 2\xi e^2 \xi e^4 + \xi e^3 \xi e^3 + 4\xi^2 e^4 + \xi e \xi e^3),$$

$$\text{tr}(\xi^2 e^7 + \xi e \xi e^6 + \xi e^2 \xi e^5 + \xi e^3 \xi e^4 + 3\xi^2 e^5 + \xi e \xi e^4), \dots$$

They are mutually commutative.

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- We consider a representation of another Lie algebra and obtain a quantum derivation/quantum argument shift operator. However, this naive operator does not satisfy the quantized argument shift method. This means that I still do not know the appropriate definition of the quantum derivation/quantum argument shift operator in the general case.
- It may be more promising to generalize this result to the general linear Lie superalgebra $\mathfrak{gl}_{m|n}$. I am currently working on this.